

Basic Phenomena of Sputter-Deposition

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Part I

Introduction, History and Aim
Generation of Atoms
Energy Transport
(low pressure case)

Part II

Transport of Atoms
Throw Distance
Gas Rarefaction
(low - high pressure transition)

A personal view: sputtering of amorphous semiconductors

(Or: How the transistor switched my life)

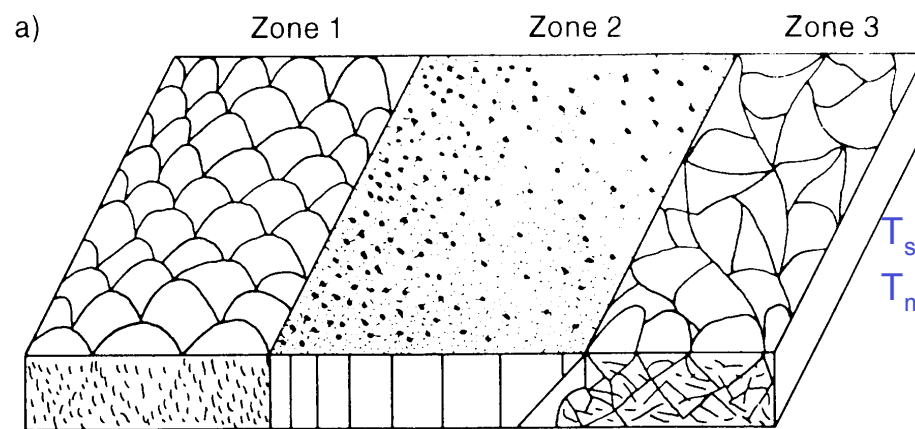
- 1947 Bardeen, Brattain and Shockley invented the transistor (1956 Nobel Prize awarded)
- 1960s Deposition of amorphous semiconductors (a-Ge, a-Si) by evaporation and sputtering aiming at thin-film electronics: Electronic defects are unsolved problem.
- 1971 William Paul (Harvard) proposed hydrogenation of a-Ge for defect passivation (namely dangling bonds) : a-Ge:H prepared by sputtering
- 1975 Spear and LeComber (Dundee) demonstrated n- and p-type doping of amorphous hydrogenated silicon (a-Si:H) by PECVD
- 1980s Intense research on a-Si:H - fundamentals of preparation and characterization, application to thin film transistors, solar cells and electro-photography (Xerography)
- 1990s Preparation of a-Si / metal (SiMal) layered structures

Why investigating energy fluxes connected to sputter-deposition?

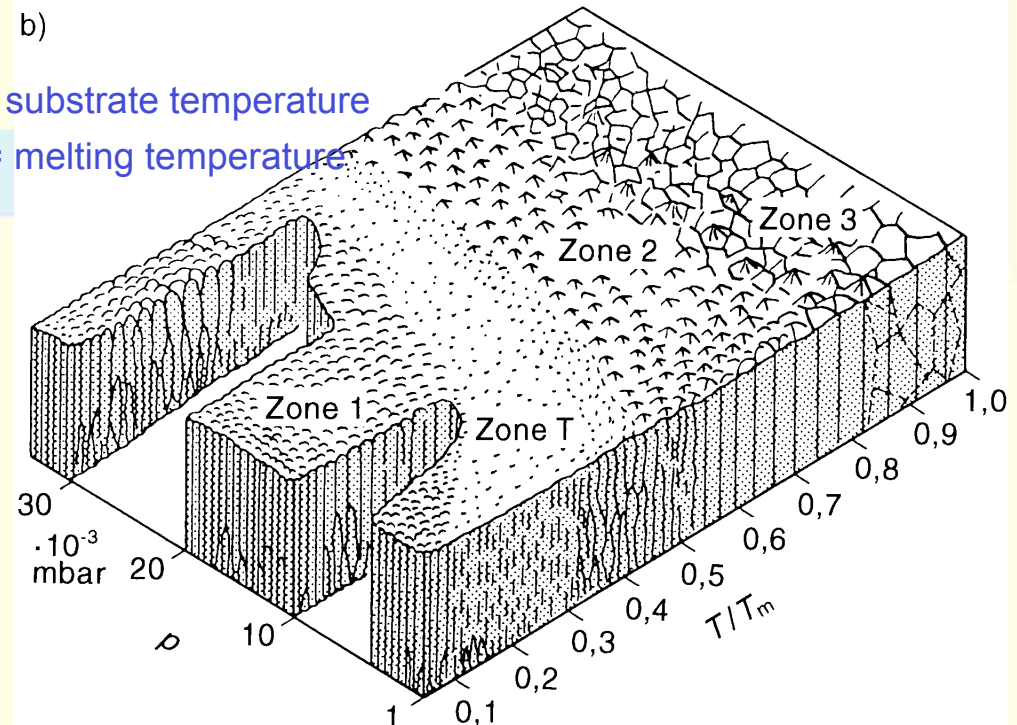
Knowledge about the energy of sputtered atoms (and other species) is useful:

- Deposition of “zero stress” films of refractory metals , which are used as conductor lines in highly integrated electronics.
- Preparation of smooth surfaces and interfaces for examples in multi-layer films for x-ray mirrors.
- Ability to prepare dense or nano-porous metal films (fractal dimension of $D_m=2.4$), which can act as field emitters or catalysts.
- Enhancement or suppression of crystallite growth and the formation of specific type of texture (for example, c-axis orientation of large AlN grains).

Structure-zone model for sputtered films



a) Evaporated films (Movchan and Demchishin)
b) Sputtered films (Thornton)



T_s = substrate temperature
 T_m = melting temperature

$T/T_m =$

Metalle < 0,26	0,26 ... 0,45	> 0,45
Oxide < 0,30	0,30 ... 0,45	> 0,45

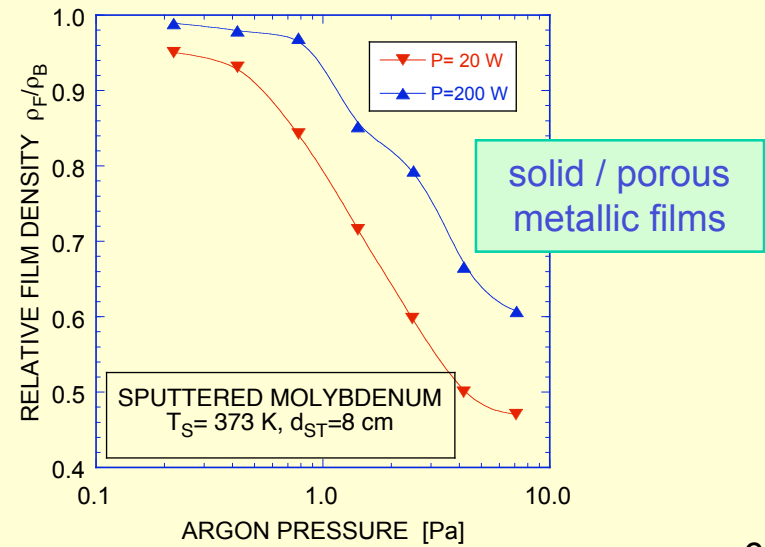
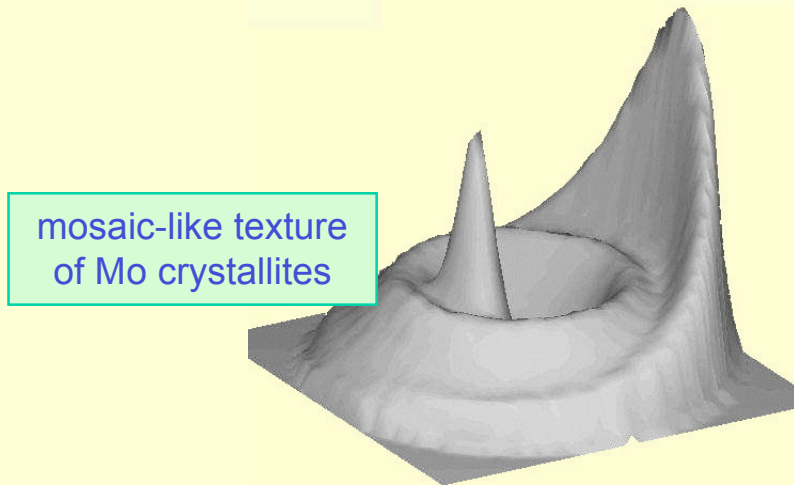
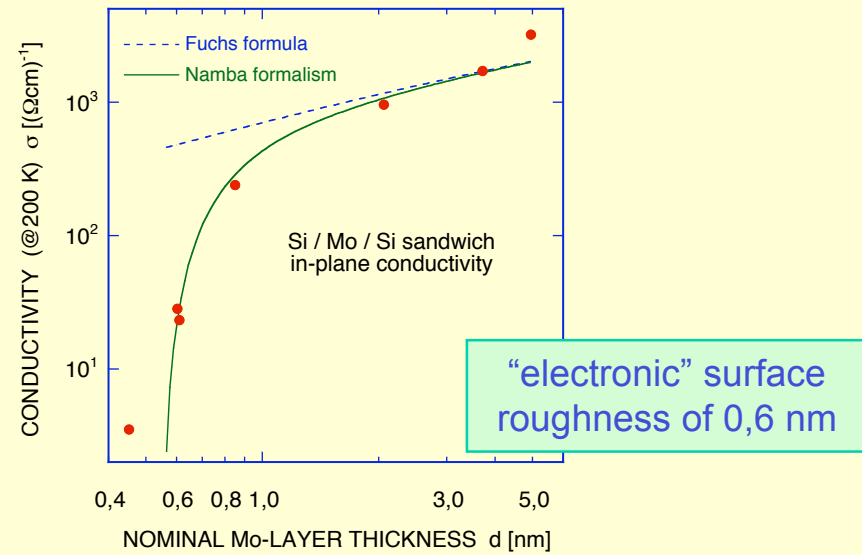
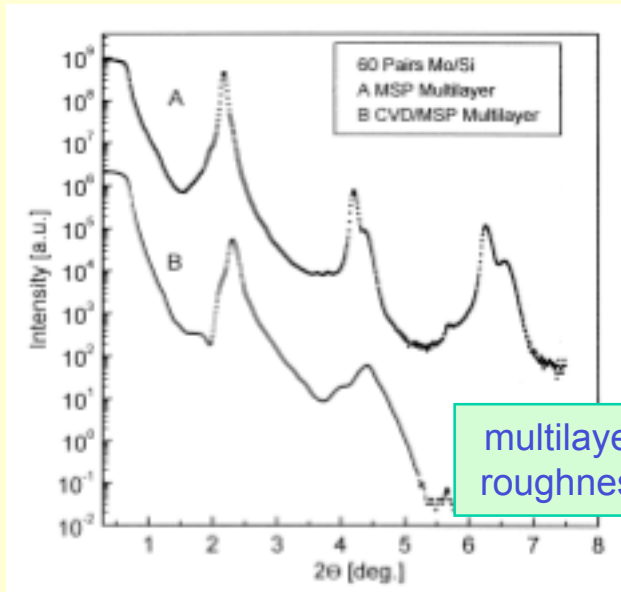
Z1: $T_s/T_m < 0.3$ no surface diffusion, initial columns of some 10 nm diam. + voids, amorphous or nano-crystalline films, columns grow to μm diam. at surface

ZT - transition zone: weak surface diffusion due to bombardment preventing microstructure

Z2: $T_s/T_m > 0.3$ strong surface diffusion, columnar crystallites of μm diam. increasing with film thickness and temperature

Z3: $T_s/T_m > 0.5$ larger non-columnar crystallites grow, smooth surfaces with grooves at boundaries

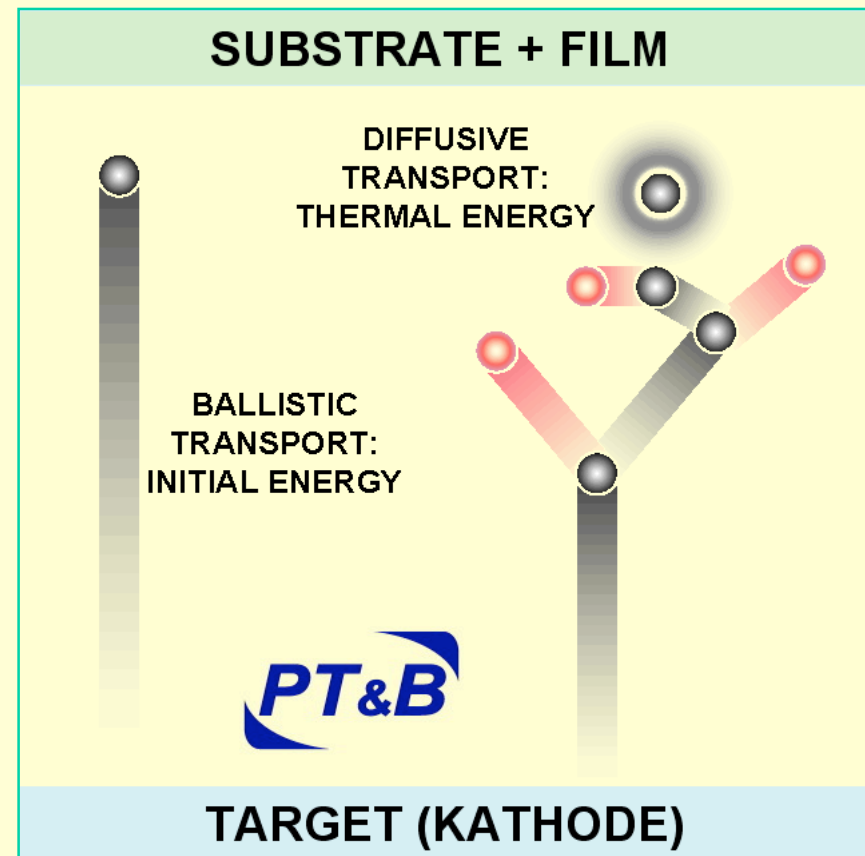
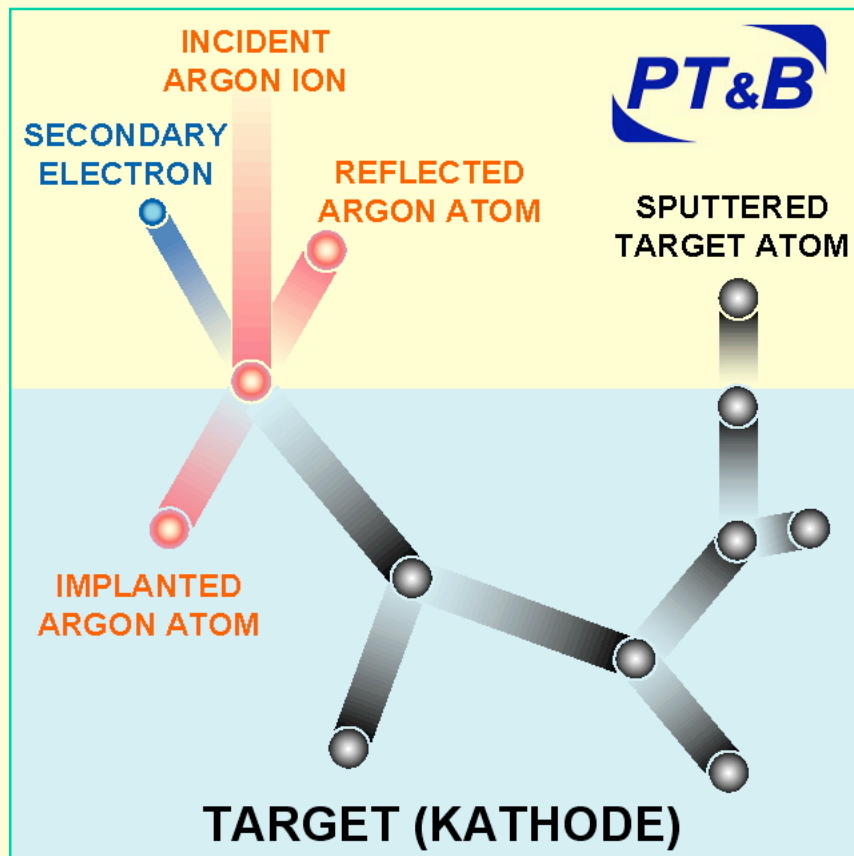
Sputtering causes novel effects on thin films !



The two processes involved: sputtering and transport

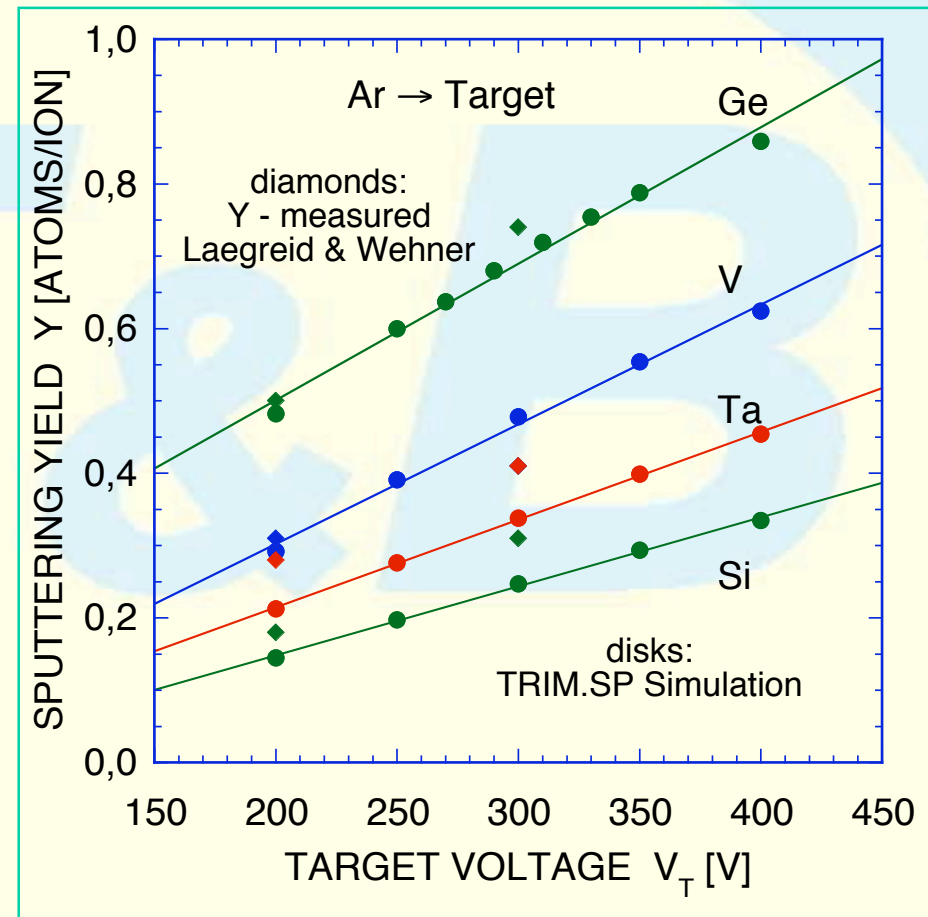
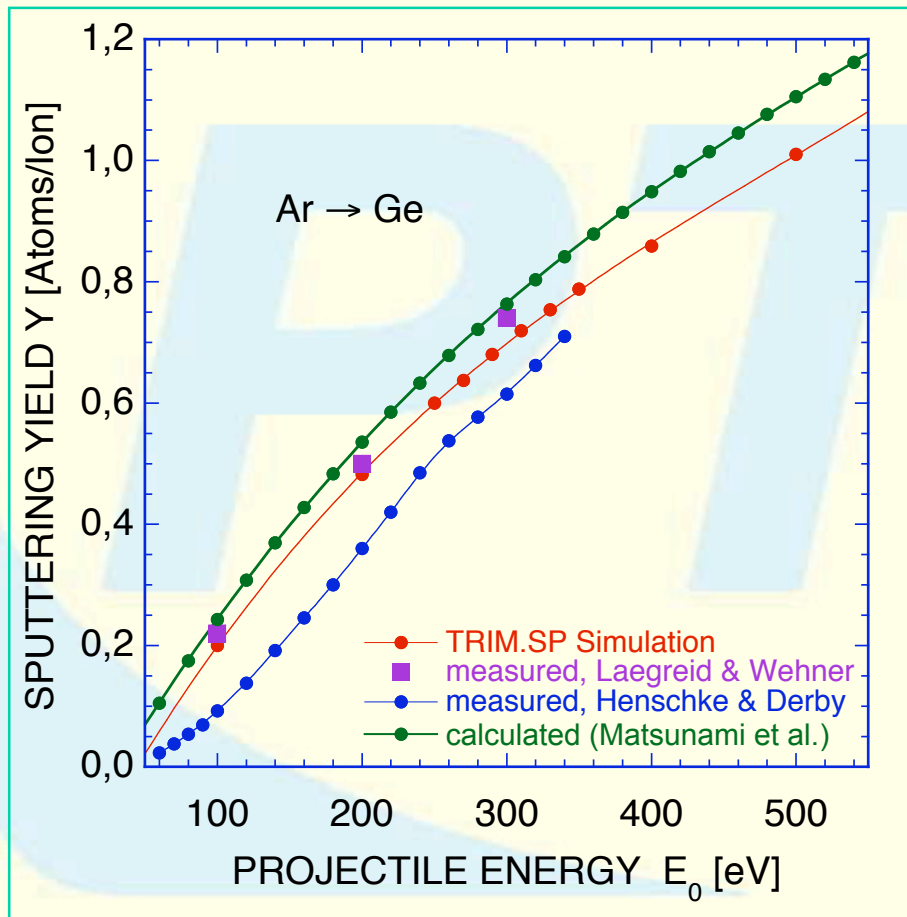
- Ar ions recombine, emission of secondary (Auger) electrons, reflection or penetration of Ar atoms
- Evolution of collision cascade of target ions
- Momentum “reversion”, energy transferred to surface atom above surface binding energy → ejection of atom

- Collisionless (ballistic) transport: target atoms conserve full initial kinetic energy (low pressure-distance product)
- Diffusive transport: multiple collisions of sputtered atoms result in complete thermalization (high pressure-distance product)



The sputtering yield Y

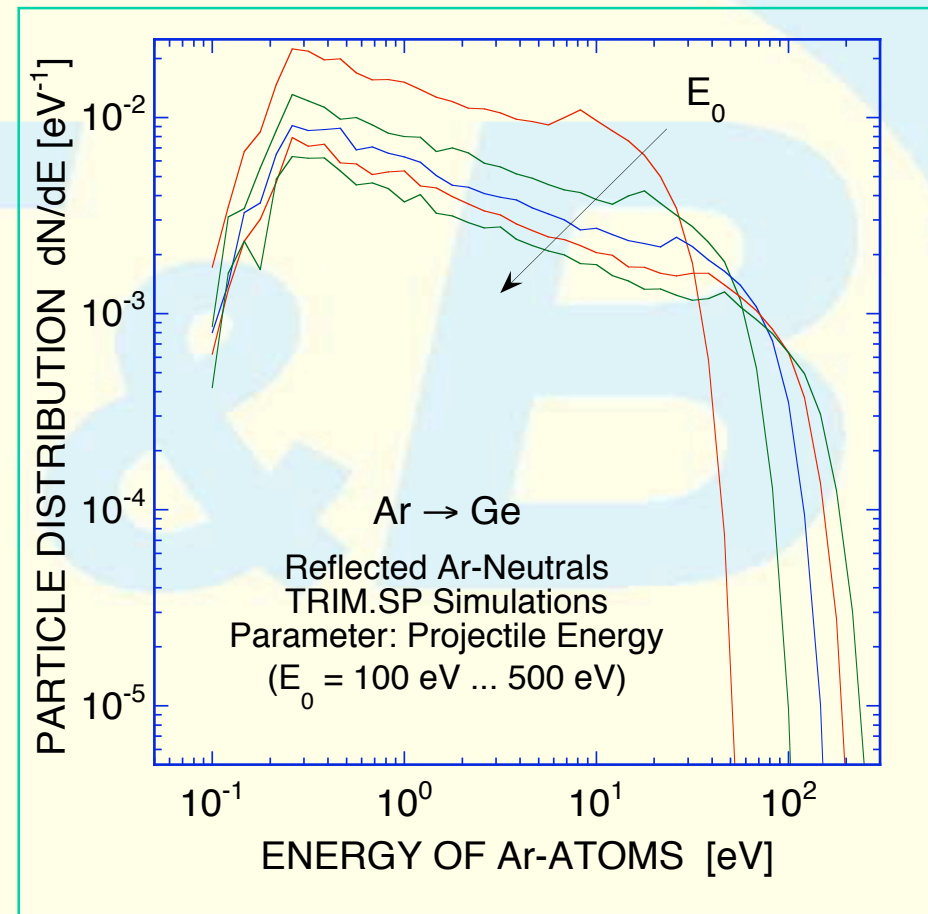
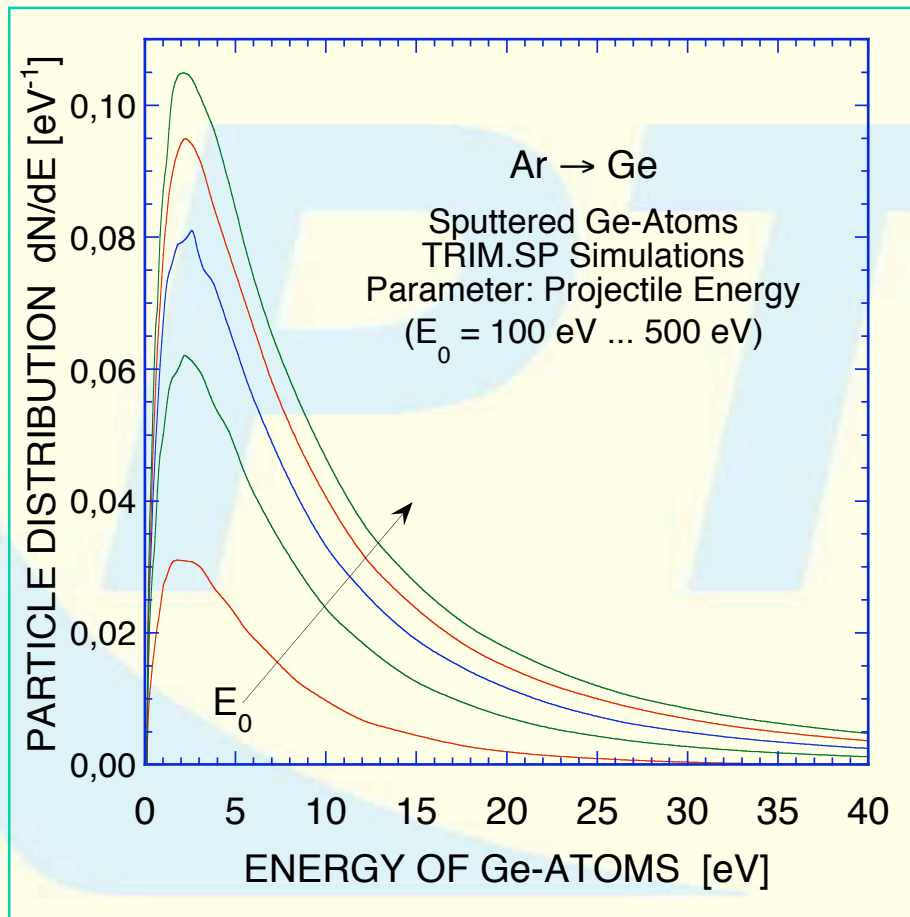
Most “popular” quantity of sputter deposition - $Y(E) =$ sputter-ejected atoms per incident ion: Depends on target and projectile mass, kinetic and surface binding energy, crystallite orientation, incident angle, roughness AND ...



Energy of sputtered and reflected atoms

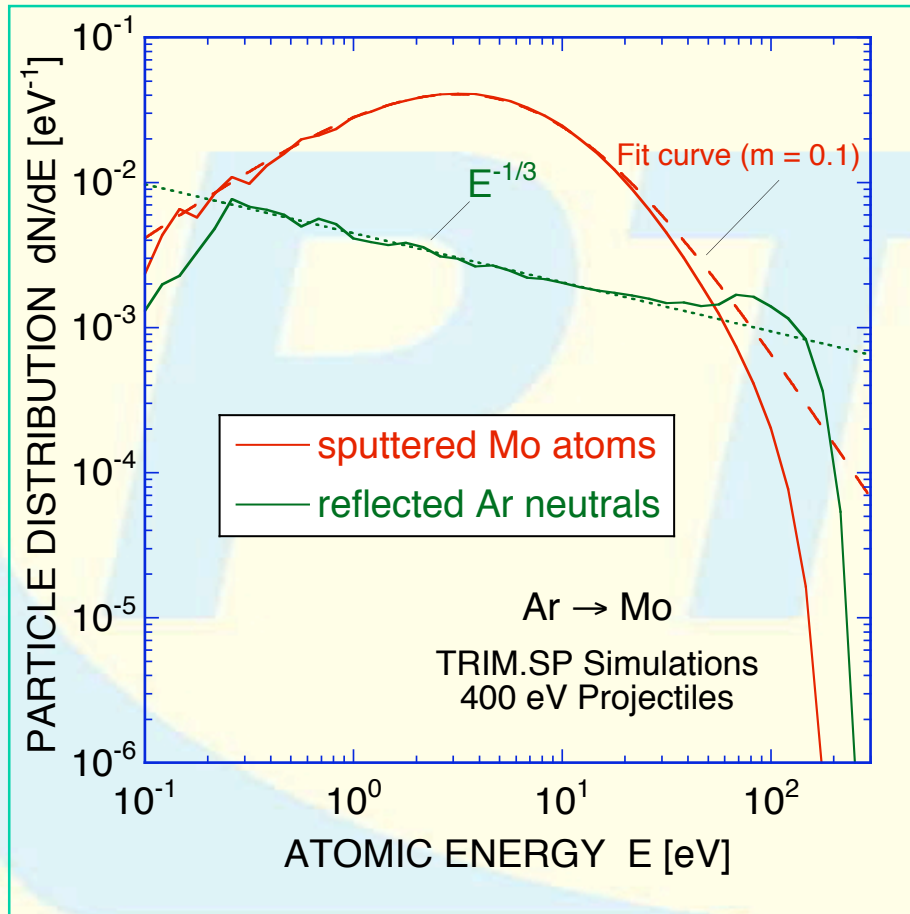
$$f(E) = \frac{E}{(E + U_0)^{3+2m}}$$

With $m=0$ Thompson distribution, maximum @ $U_0/2$
 U_0 - surface binding energy

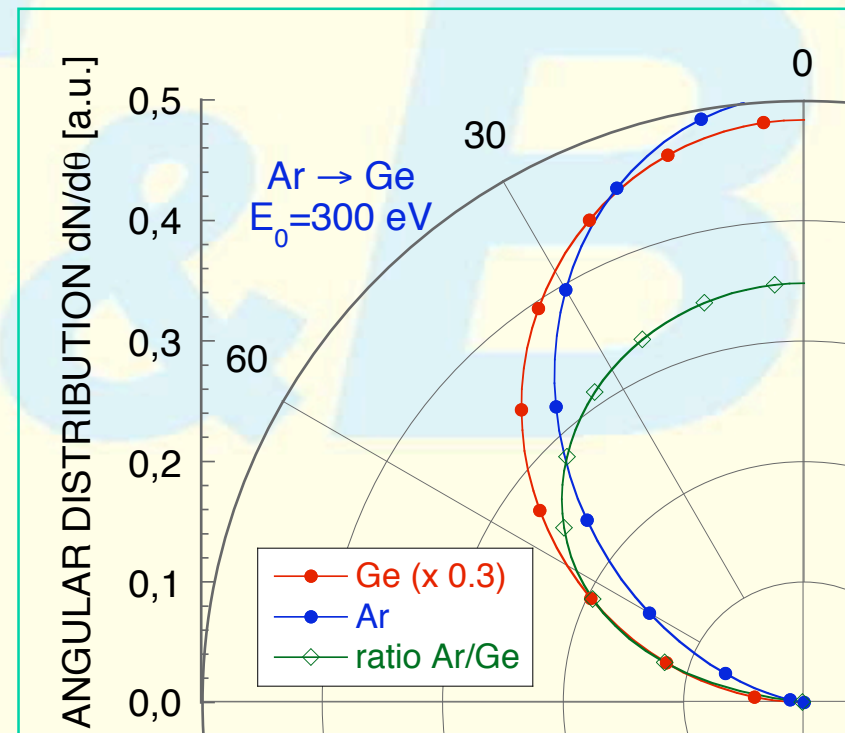


Energetic and angular distribution of sputtered and reflected atoms

Cut-off energy is about E_{\max} !



$$dN/d\theta = \cos^{\nu} \theta, \quad \nu \approx 1$$



Estimate of energy input per atom by elemental sputter-deposition

$$\langle E_t \rangle = I_{tot} / \Phi_{at}$$

$\langle E_t \rangle$ measured by calorimetric method (power density I_{tot}) and atomic deposition rate Φ_{at}

$$\langle E_t \rangle = U_0 + \langle E_{at} \rangle + \langle E_{Ar} \rangle / at + \langle E_{Plasma} \rangle / at$$

$\langle E_t \rangle$ calculated from quantities resulting from TRIM.SP forward simulation
 U_0 - surface binding energy

$\langle E_{at} \rangle$ - average energy of sputtered atoms

$$\langle E_{at} \rangle = \frac{\int_0^{E_{max}} f(E) E dE}{\int_0^{E_{max}} f(E) dE}$$

$$E_{max} = k \frac{4M_I M_T}{(M_I + M_T)^2} E_0 - U_0$$

with E_{max} maximum transferred energy and $k = f(M) = 0.1 \dots 0.4$

$\langle E_{at} \rangle$ - analytical approach

$$\langle E_{at} \rangle = U_0^{2/3} E_{max}^{1/3}$$

$\langle E_{Ar} \rangle / at$ - average energy of reflected argon per deposited atom
 R_N, R_E - Particle und Energy reflection coefficient

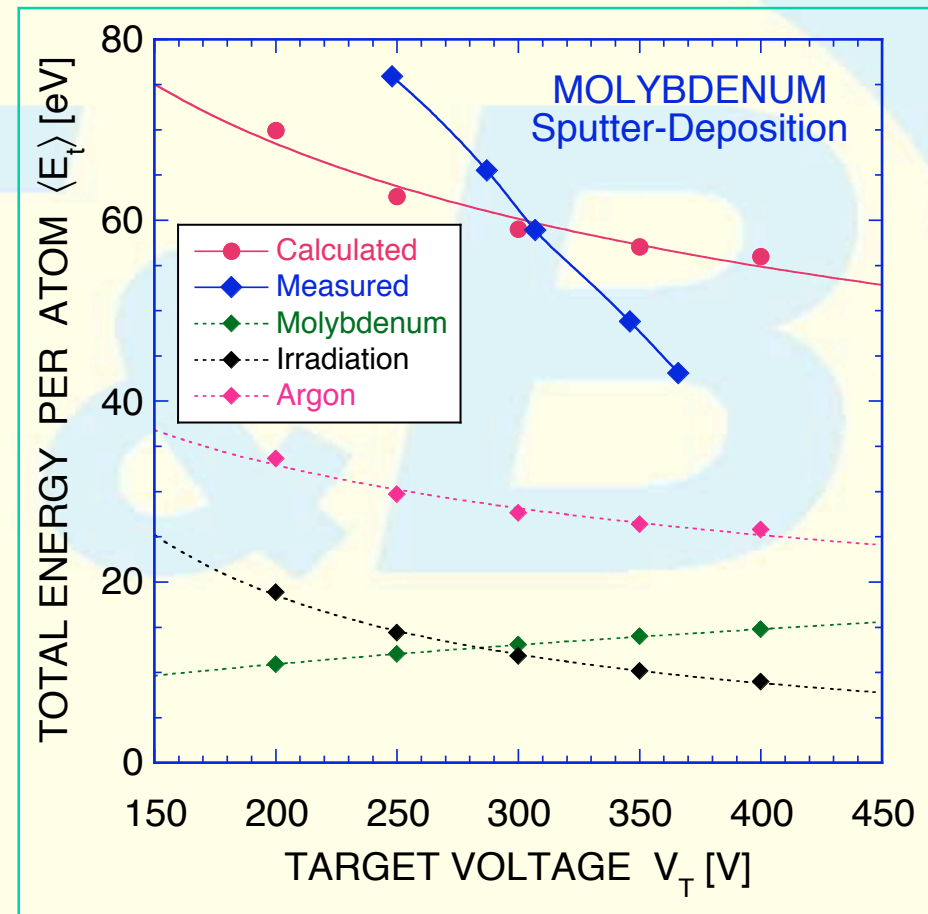
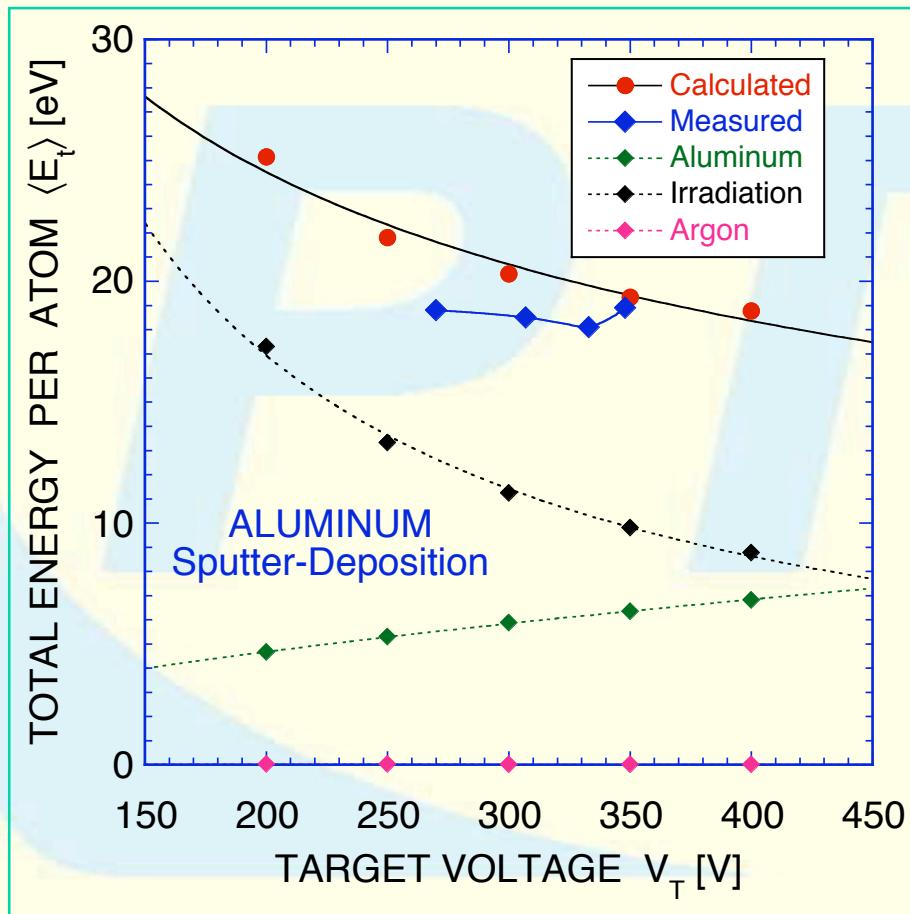
$$\langle E_{Ar} \rangle / at = E_0 \frac{R_E}{R_N} \frac{R_N}{Y}$$

$\langle E_{Plasma} \rangle / at$ - average plasma irradiation energy per deposited atom
 Y - sputtering yield

$$\langle E_{Plasma} \rangle / at = \frac{5,33eV}{Y}$$

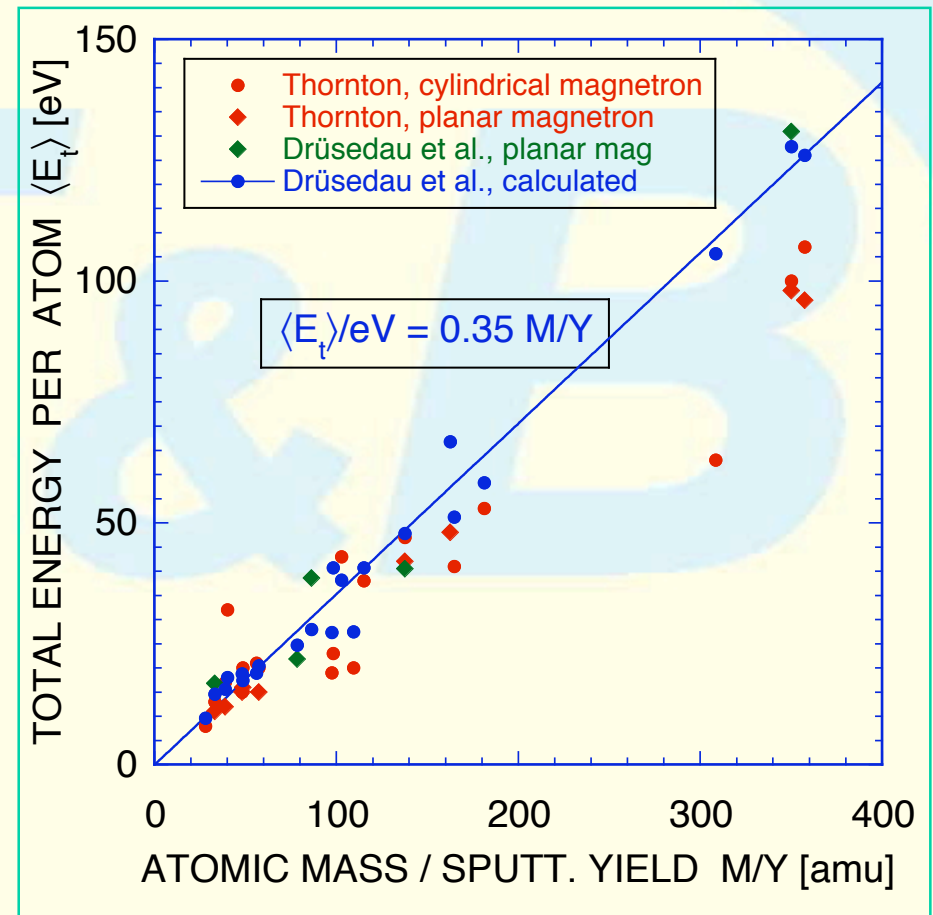
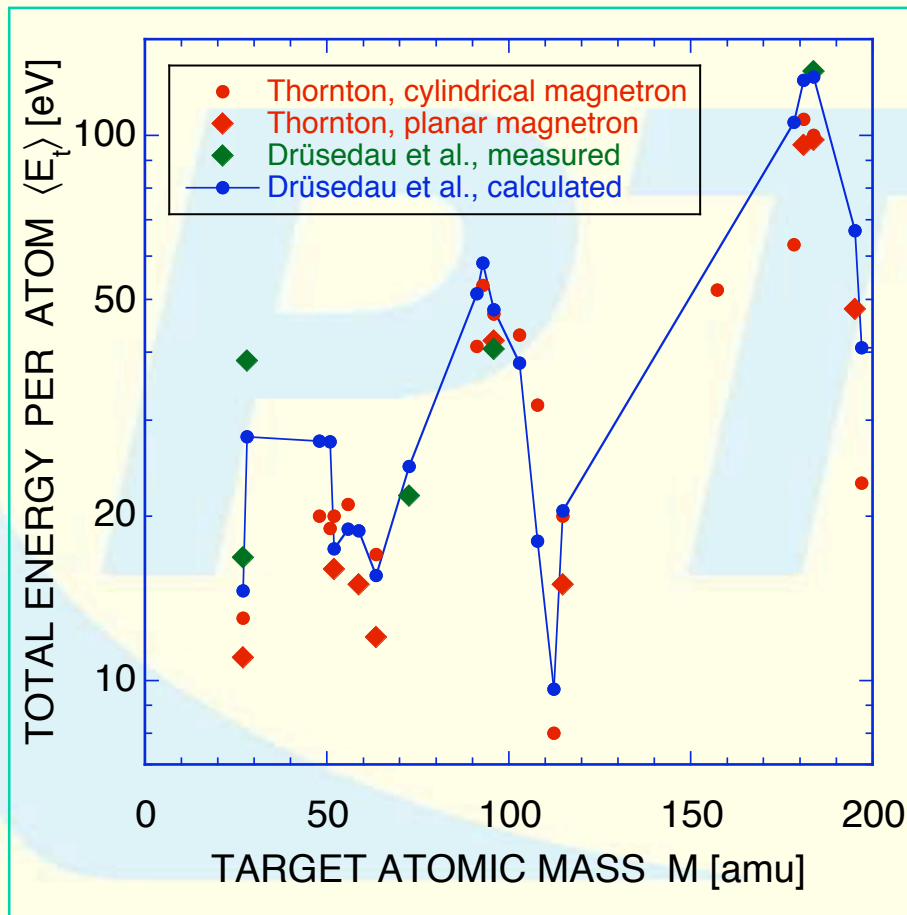
Results: Energy input by elemental sputter-deposition

Comparison of values measured to the results of forward simulation (with the components due to kinetic energy of the sputtered atoms and reflected argon neutrals and plasma irradiation)



Results: Energy input by elemental sputter-deposition

Comparison of values measured to the results of forward simulation. The dependence of $\langle E_t \rangle$ on the atomic mass (left) is superimposed by the surface binding energy (sublimation enthalpy) of the elements.



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Do YOU know the meaning of the word “sputter” ?

“What is all that sputtering nonsense anyway?”
(from “Glow Discharge Processes” by Brian Chapman)

- a medical and a sputtering conference took place at Imperial College
- the well-known scientific phenomenon: conference system tends towards a condition of being in the bar
- a well-oiled medic demanded to know: “What is all that ...”
- reply: “Well, we’re in a branch of medical profession, too - in speech therapy actually. Sputtering is like stuttering, except, our chaps say p...p...p instead of t...t...t
- the medic warmly thanked his newly discovered colleague

“sputter” appeared in English 1598, adapted from the the words “sputteren” (Dutch) and “sputterje” (West Frisian)

English Dictionary (The Shorter OED 1957):

“To spit out in small particles and with a characteristic explosive sound ... His tongue was too large for his mouth; he stuttered and sputtered (1878)”

Experimental

Sputter-deposition

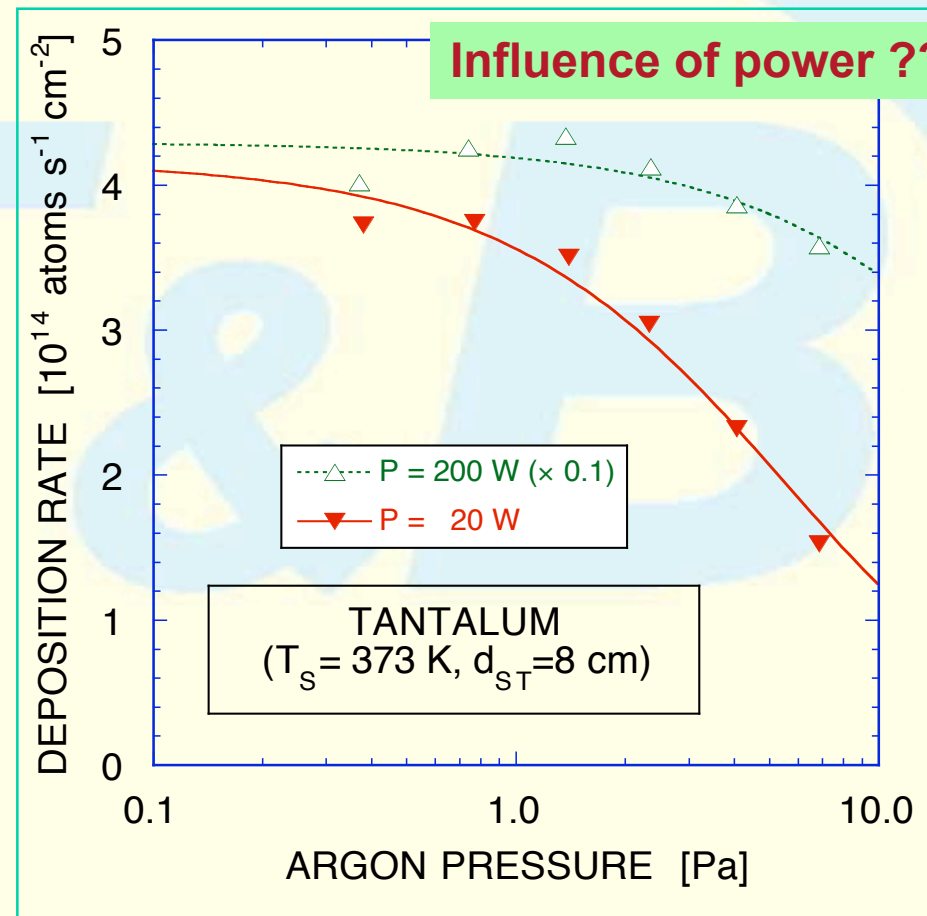
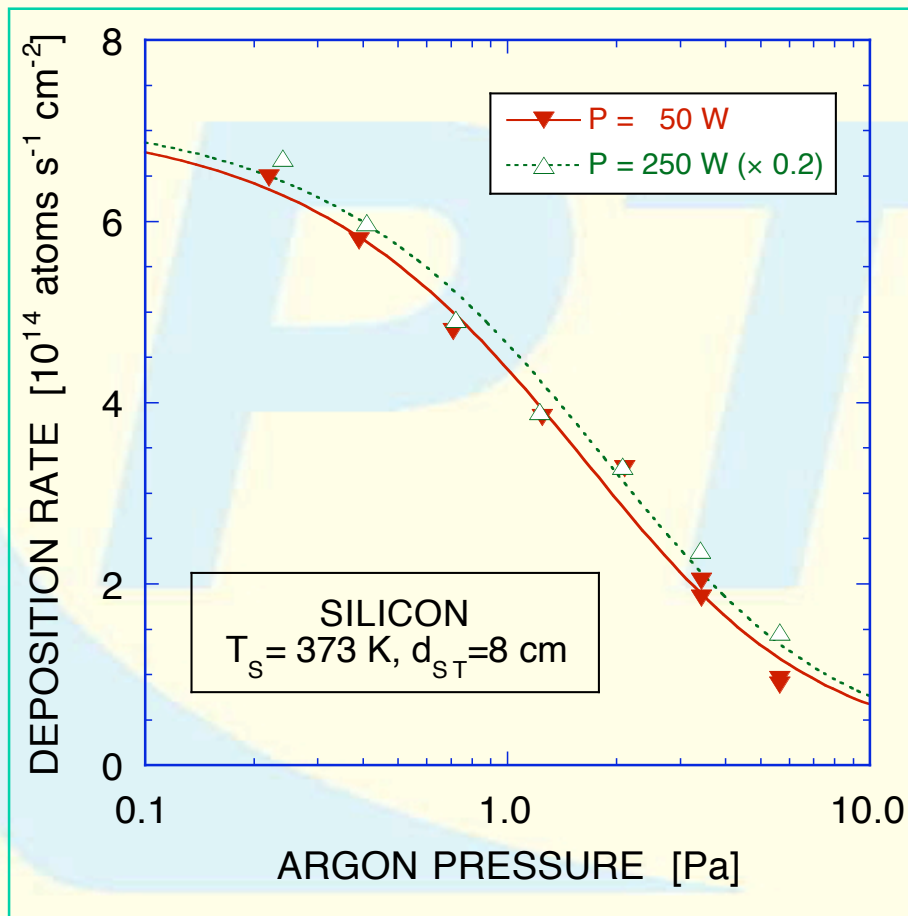
- 90 mm diameter Targets: Al, Si, Ti, V, Cr, Ge, Mo, Ta, W
- Sample sets: variation of pressure $p_{Ar}=0.2 - 7.0$ Pa
- Power (typical): $P_{DC}=20$ W (low power – 0.3 Wcm⁻²) **and**
 $P_{DC}=200$ W (high power – 3.1 Wcm⁻²)
- Substrate-to-target distance of $d_{ST}=80$ mm
- Substrate Temperature of $T_s=373$ K
- Film thickness of 100 nm and 500 nm, respectively
- Calorimetric measurements due to sample heating and cooling
- Gas density reduction measurements

Film characterization

- Film thickness: Talystep instrument and X-ray reflectivity
- Specific gravity: X-ray reflectivity
- Determination of atomic deposition assuming unity sticking probability

Pressure-dependent atomic deposition rates

Throw distance: characteristic decay length of deposition rate as a function of distance from target at given pressure → more general: consideration of pressure - distance product !



Pressure-dependent flux of sputtered atoms

One-dimensional model by Keller and Simmons

Flux of ballistic atoms Φ_B at substrate:

$$\Phi_B(pd_{ST}) = \Phi_0 \exp\left(-\frac{pd_{ST}}{p_0 d_0}\right)$$

Φ_0 “zero-pressure” flux
 d_{ST} substrate-to-target distance
 $p_0 d_0$ characteristic pressure-distance (throw distance)

Single-step thermalization of ballistic atoms

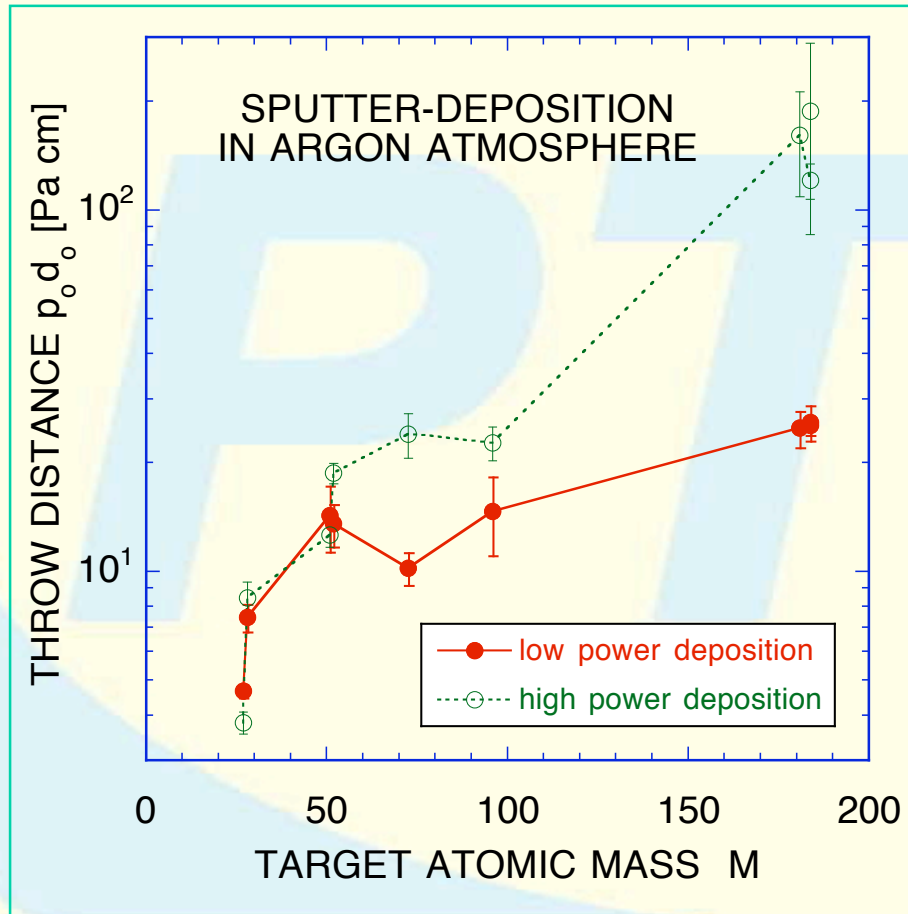
→ source of diffusing thermal atoms of flux Φ_D at substrate:

$$\Phi_D(pd_{ST}) = \Phi_0 \left\{ \frac{p_0 d_0}{pd_{ST}} - \left(1 + \frac{p_0 d_0}{pd_{ST}}\right) \times \exp\left(-\frac{pd_{ST}}{p_0 d_0}\right) \right\}$$

Total atomic flux Φ at substrate:

$$\Phi(pd_{ST}) = \Phi_B + \Phi_D = \Phi_0 \frac{p_0 d_0}{pd_{ST}} \left\{ 1 - \exp\left(-\frac{pd_{ST}}{p_0 d_0}\right) \right\}$$

Throw distances are influenced by multiple parameters



Atomic mass:

Influence on throw distance is obviously related to fundamental atomic processes

Sputtering power:

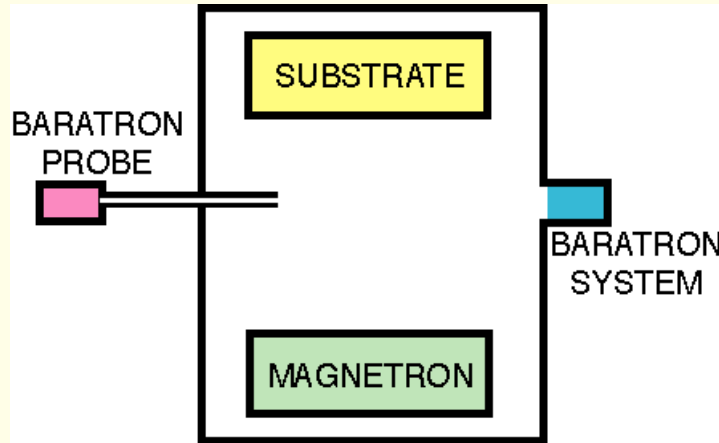
Influence on throw distance is related to process (apparatus) specific effects

The problem:

To identify and to separate the process-specific parameters from the effects measured to obtain general laws.

→ Investigation of gas rarefaction

Setup for measuring gas rarefaction (“sputtering wind”)



sputtering chamber and manometer at the tube’s end
 a system of two connected vessels each characterized
 by a temperature and pressure of p_S, T_S and p_P, T_P .

transition range between molecular and
 viscous flow: the ratio of the pressure values
 at both ends of the tube is described by:

$$\frac{p_P}{p_S} = \frac{f(X) + \sqrt{T_P/T_S}}{f(X) + 1}, X = d_T p_P$$

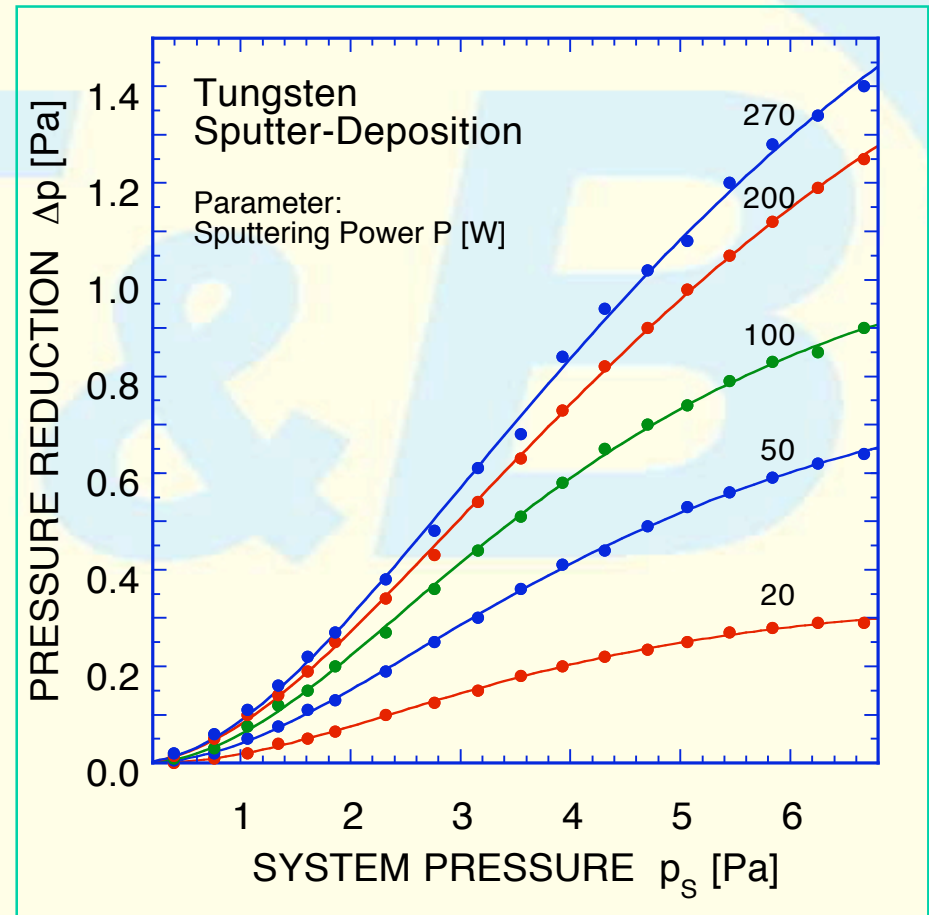
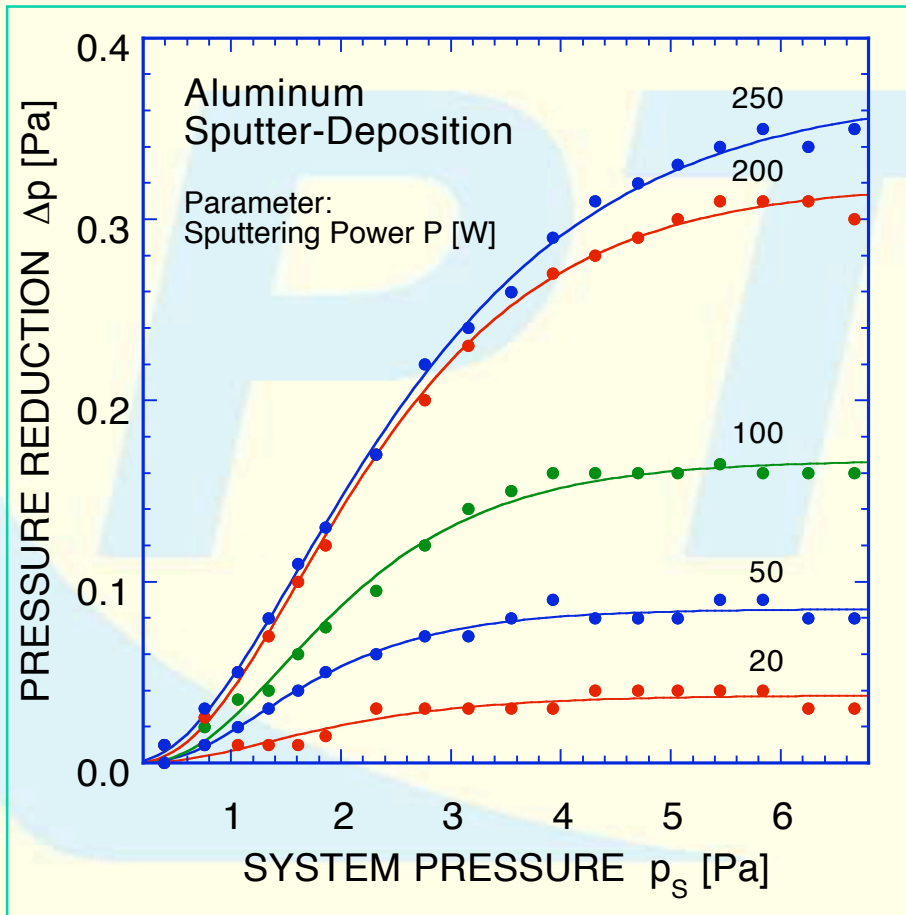
$$f(X) = A(X/T^*)^2 + B(X/T^*) + C(X/T^*)^{1/2}, T^* = (T_S + T_P)/2$$

Possible expression of $f(x)$
 with empirical constants
 A, B and C

Raw data: Probe pressure is reduced to system pressure

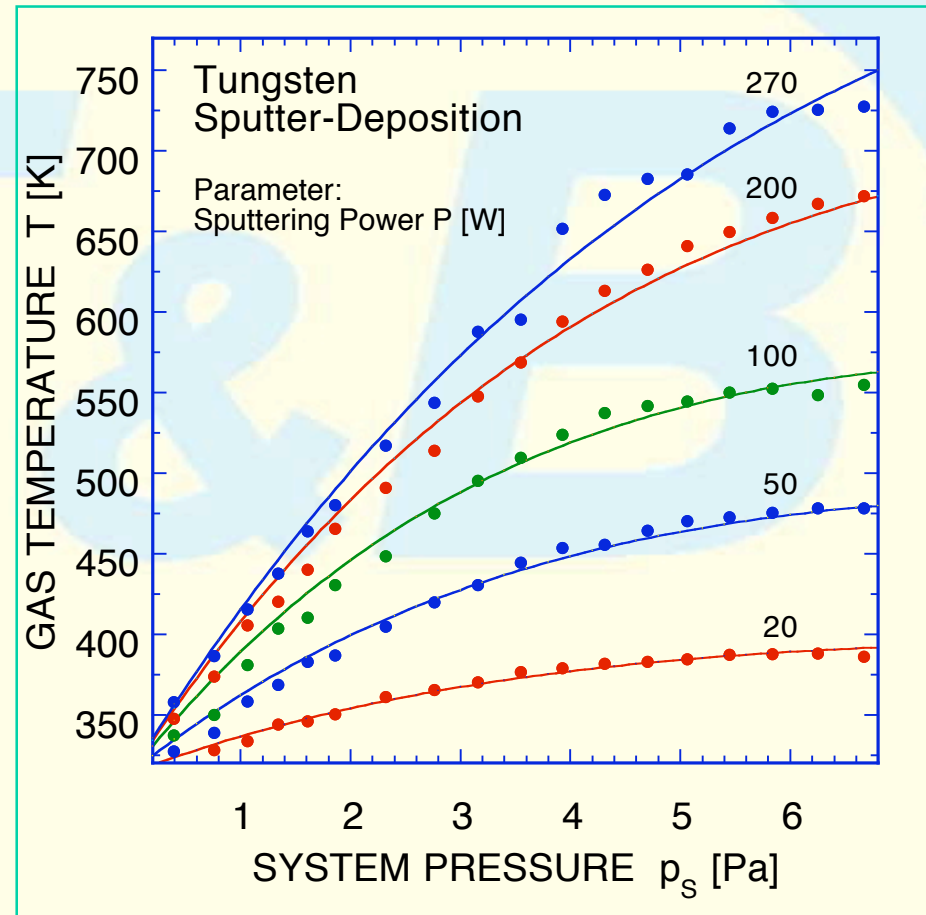
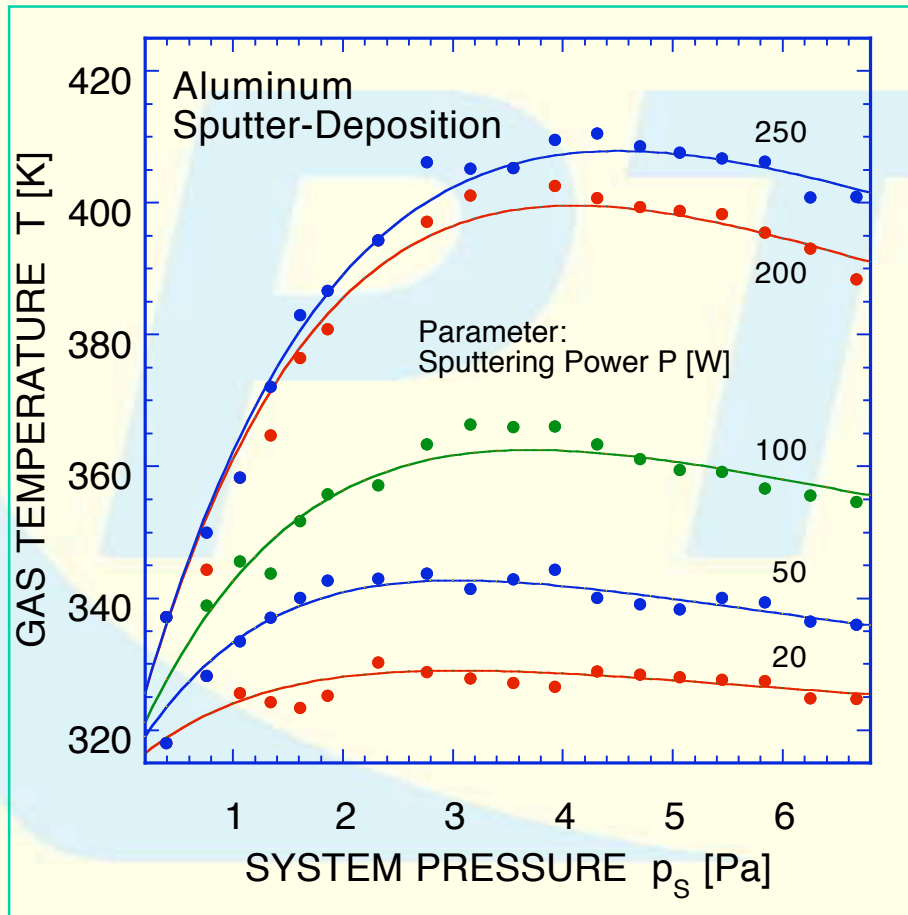
empirical fit curve:

$$\Delta p = p_S - p_P = m_1 \{1 - \exp(-m_2 p_S)\}^{m_3}$$



Gas temperature as a function of pressure (sputtering power)

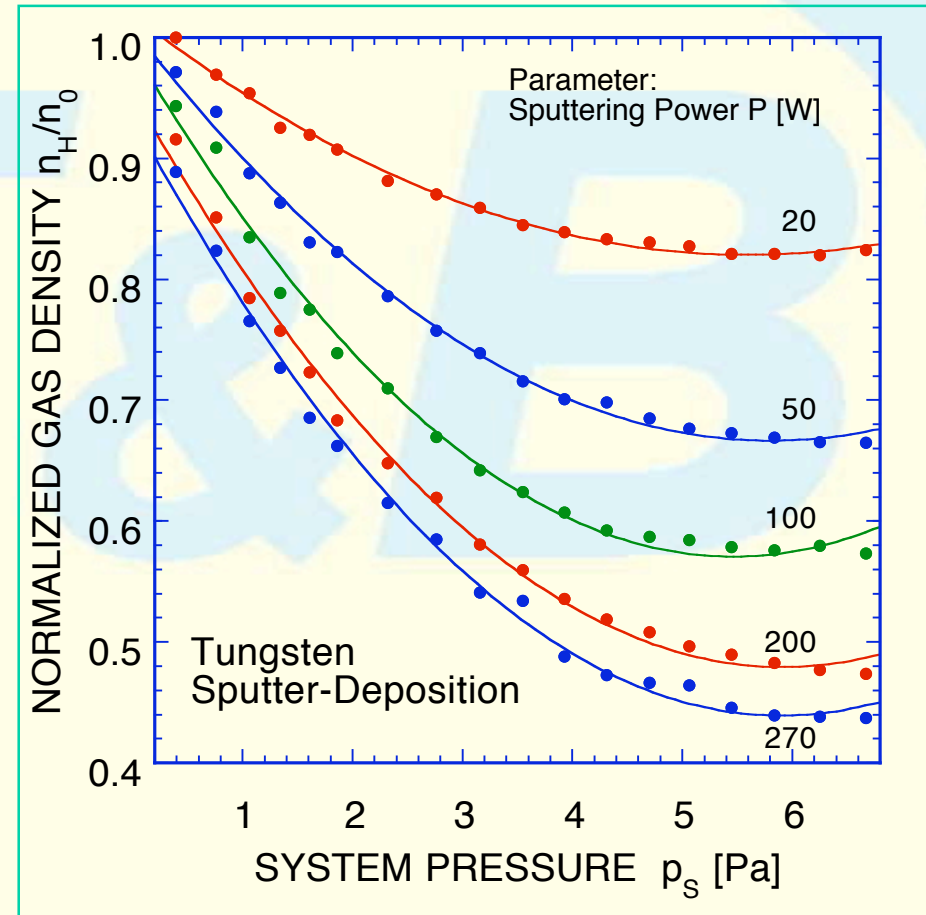
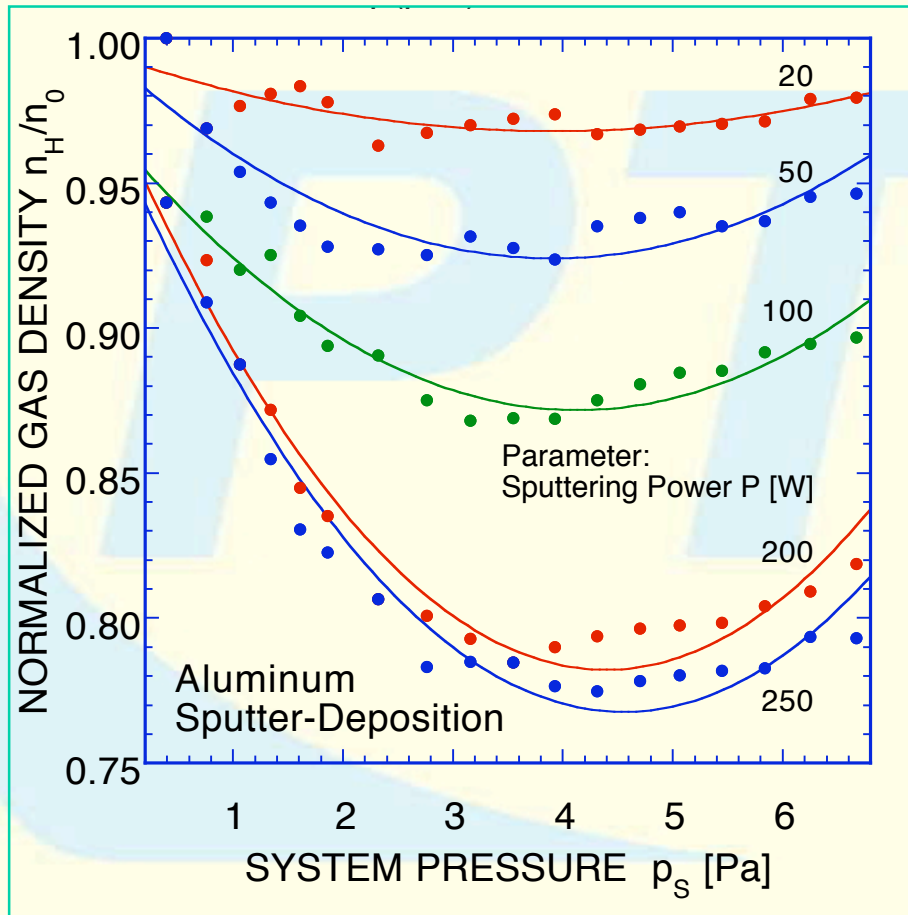
System (gas) temperature calculated from systems pressure and temperature (p_S, T_S) and probe pressure p_P . Fit curves according to the analytical model.



Gas density (normalized) as a function of pressure (power)

From temperature T_H measured, gas density n_H in the heated plasma region is determined (n_0 and T_0 gas density, temperature in the “cold”)

$$\frac{n_H}{n_0} = \frac{T_0}{T_H}$$



One-dimensional model for gas heating

In a sheath of dx around a position x a heating power $P(x)$ is generated according to loss in ballistic atoms with average kinetic energy $\langle E_A \rangle$

$$P(x) = -\langle E_A \rangle \frac{d\Phi}{dx} dx = \langle E_A \rangle \frac{p\Phi_0}{p_0 d_0} \exp\left(-\frac{px}{p_0 d_0}\right) dx$$

Performing the integration of the differential equation of the one-dimensional steady-state heat flow:

$$T(x) = -\langle E_A \rangle \Phi_0 \frac{p_0 d_0}{\kappa p} \exp\left(-\frac{px}{p_0 d_0}\right) + c_1 x + c_2$$

Boundary conditions:

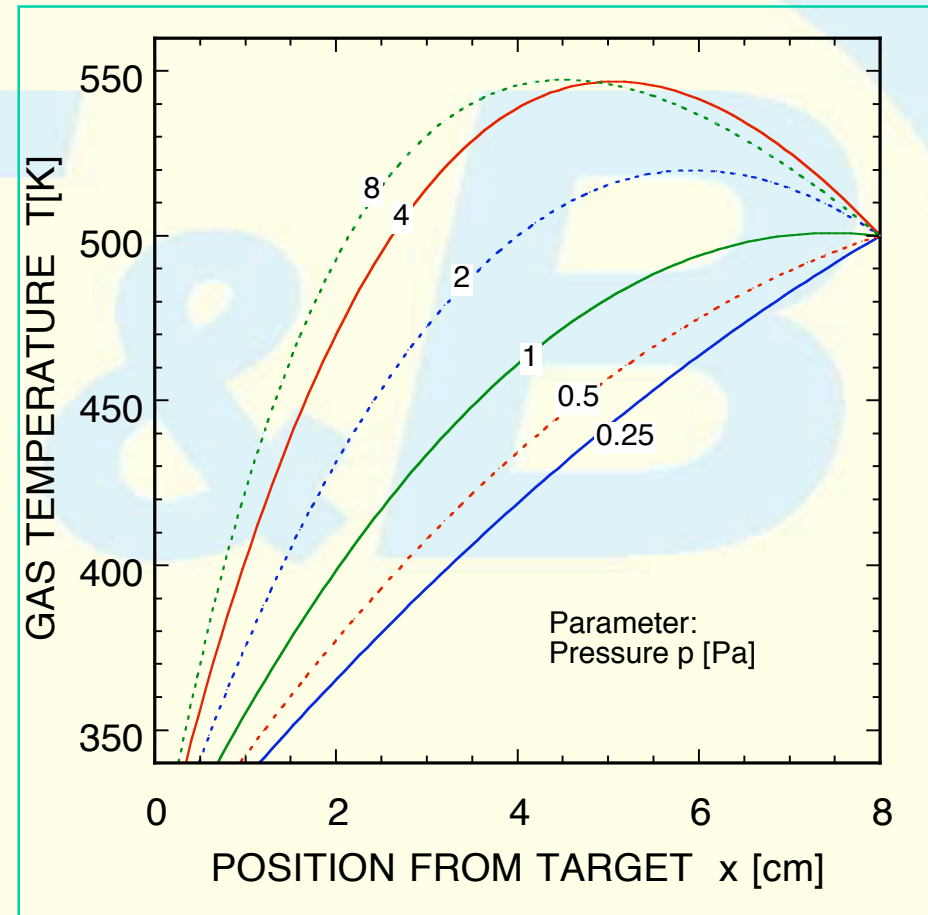
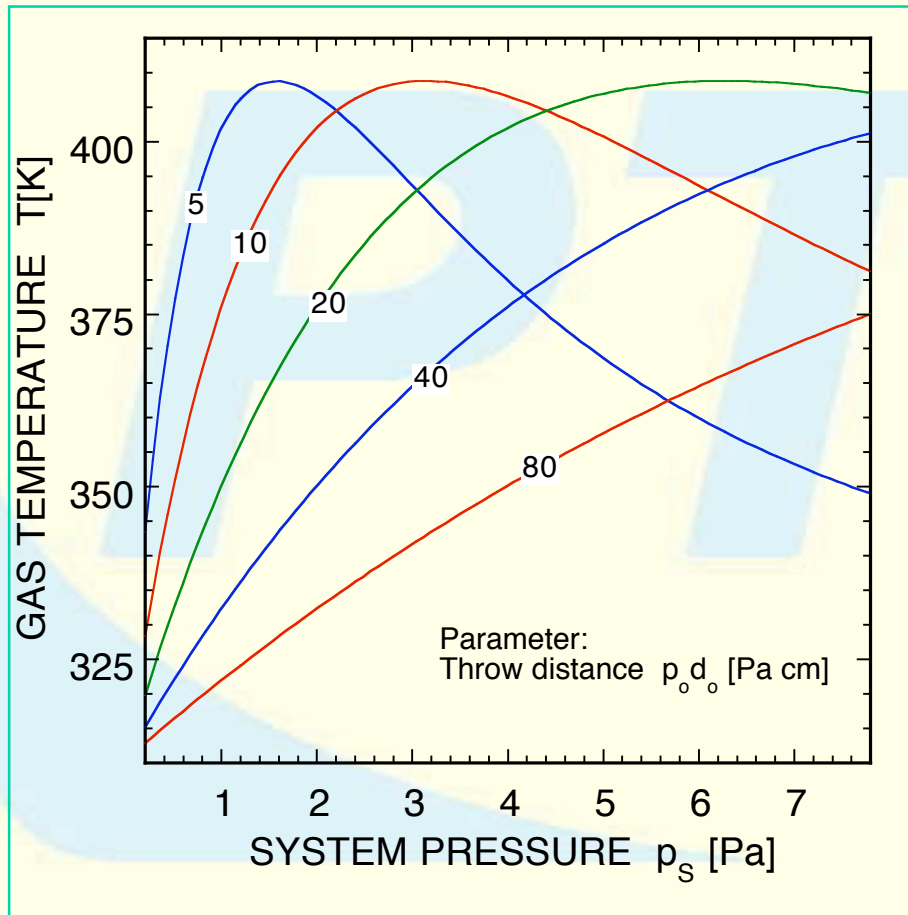
$T=T_T$ (target position, $x=0$), $T=T_S$ (substrate position $x=d_{ST}$), determination of c_2 and c_1

$$c_2 = T_T + \langle E_A \rangle \Phi_0 \frac{p_0 d_0}{\kappa p}$$

$$c_1 = \frac{1}{d_{ST}} \left\{ T_S + \langle E_A \rangle \Phi_0 \frac{p_0 d_0}{\kappa p} \exp\left(-\frac{pd_{ST}}{p_0 d_0}\right) - c_2 \right\}$$

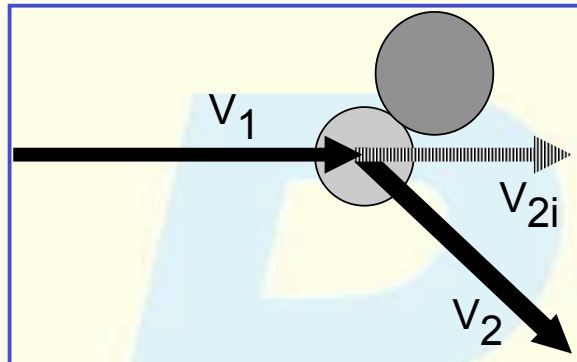
Results of model calculation

Model parameters: throw distance $p_0 d_0 = 20 \text{ Pa cm}$, average kinetic energy $\langle E_A \rangle = 10 \text{ eV}$, pressure $p = 4 \text{ Pa}$, position $x = 4 \text{ cm}$, atomic flux $\Phi_0 = 2 \cdot 10^{15} \text{ atoms s}^{-1} \text{ cm}^{-2}$.



The velocity persistence ν (Jeans 1954)

ν - a measure of momentum conservation in initial direction:

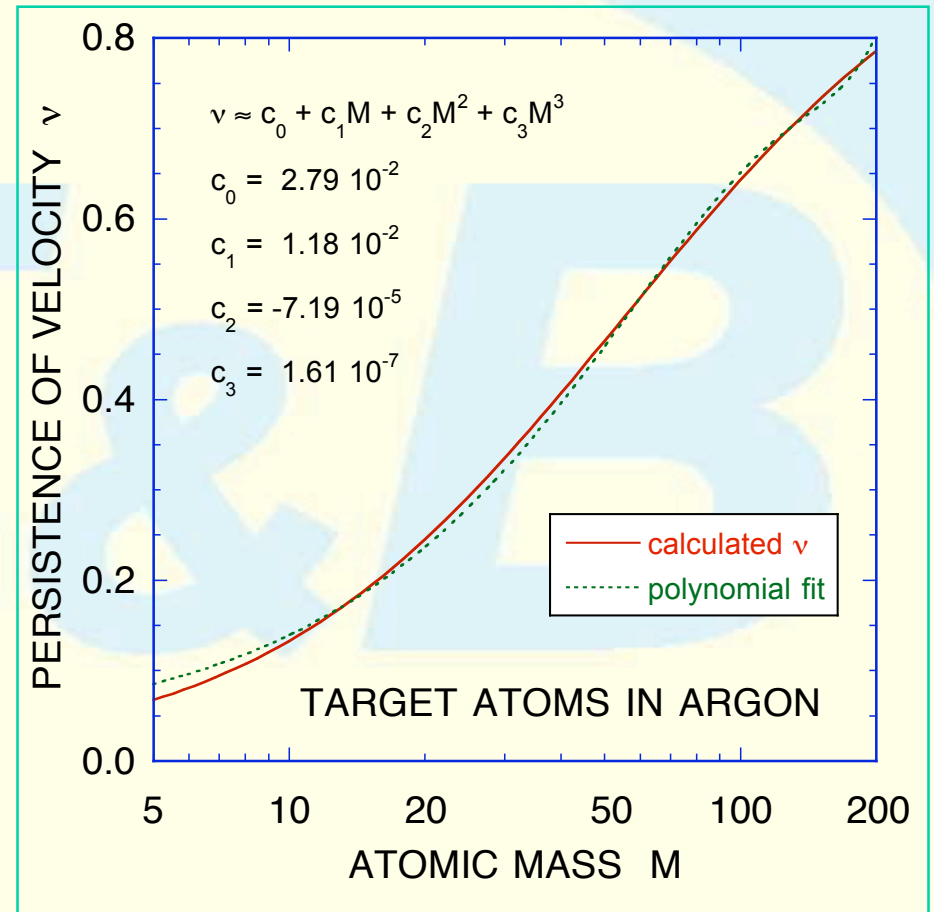


$$\nu = \frac{v_{2i}}{v_1}$$

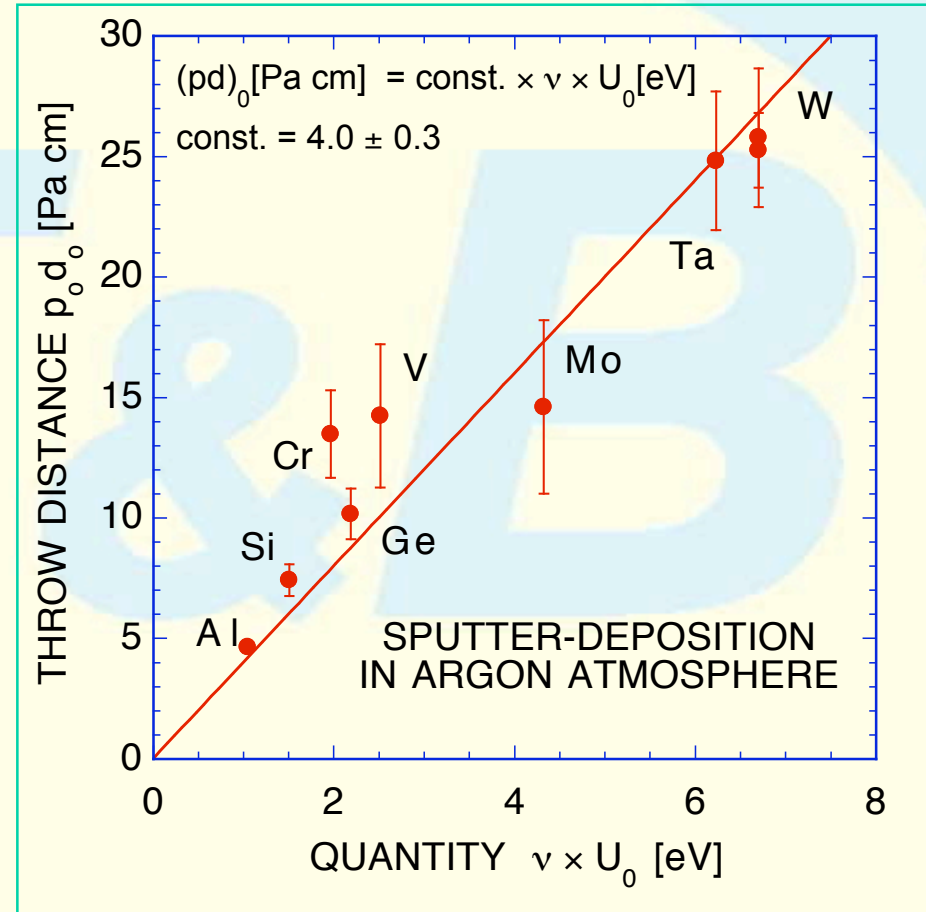
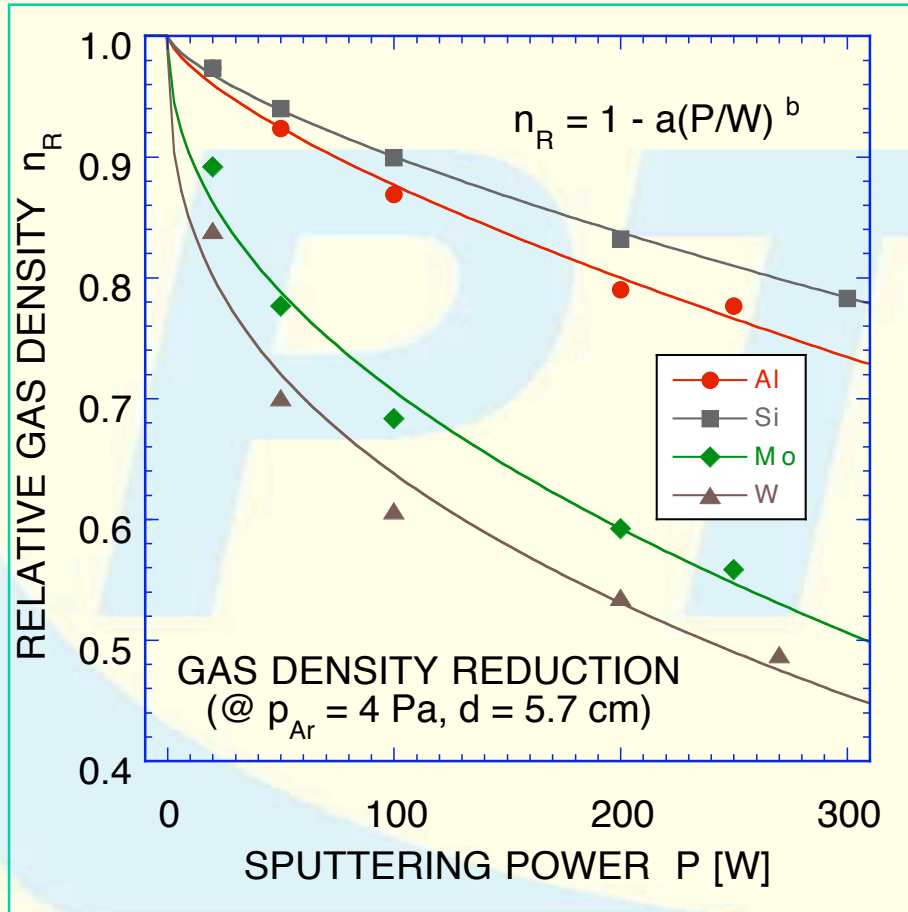
Averaging \rightarrow analytical expression for ν
(sputtered atoms of atomic mass M_S and gas M_G):

$$\nu = \frac{1 - M}{1 + M} + \frac{2M}{1 + M} v_e, \quad v_e \approx 0.4, \quad M = \frac{M_G}{M_S}$$

$$v_e = \frac{\ln[\sqrt{1 + M} + \sqrt{M}]}{4\sqrt{M^3(1 + M)}} + \frac{2M^4 + 5M^3 + 3M^2 - M - 1}{4M(1 + M)^3}$$



Avoiding gas rarefaction, there is a law for the throw distance !



A simple empirical law for the throw distance !

$$p_o d_o [Pacm] = (4.0 \pm 0.3) \nu U_o [eV]$$

fairly equivalent to:

$$p_o d_o [Pacm] = \{-1.9 + 0.96 \ln(M/amu)\} U_o [eV]$$

The heavier and faster sputtered atoms are, the farer they fly!

(That was, what we already felt 😊😐😞😐😄)

Thank you for your patience!